

# Testing Single RPV Squark Production at the LHC

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*work in progress, arXiv:151x.yyyyy*

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# Status of Natural (vanilla) SUSY

At the end of Run I,

- $m_{\tilde{g}} \gtrsim 1.4 \text{ TeV}$
- $m_{\tilde{q}_{1,2}} \gtrsim 1 - 1.3 \text{ TeV}$
- $m_{\tilde{t}} \gtrsim 700 - 800 \text{ GeV}$
- $m_H = 125 \text{ GeV} \xrightarrow{?} \tilde{m} \sim 1 - 10 \text{ TeV}$

Upper limits on fine-tuning at least at the % level.

This has lead to renewed interest on alternatives, such as  $R$ -parity violation (RPV), on both the **experimental** and **theory** side.

$$W_{RPV} \sim \mu'_i L_i H_u + \lambda_{ijk} L_i L_j \ell_k^c + \lambda'_{ijk} L_i Q_j d_k^c + \lambda''_{ijk} u_i^c d_j^c d_k^c$$

$B$  and  $L$  violation, leading to rapid proton decay, e.g.

$$\tau_p > 10^{34} \text{ years} \implies |\lambda'_{11k} \lambda''_{11k}| \lesssim 2 \times 10^{-27} \left( \frac{m_{\tilde{d}_{kR}}}{100 \text{ GeV}} \right)^2$$

A  $Z_2$  symmetry,  $R_p = (-1)^{2S+3B+L}$  removes the problem.

Also stable LSP (DM), MET at colliders.

# Baryonic RPV

Still, proton is stable is  $B$  or  $L$  separately conserved(\*). I will consider baryonic RPV,

$$W = \frac{\lambda''_{ijk}}{2} u_i^c d_j^c d_k^c, \quad i, j, k = 1, 2, 3, j \neq k$$

Flavor physics constraints ( $\Delta B = 2$  transitions,  $K - \bar{K}$ ,  $n - \bar{n}$ ):

[Barbier et al., [hep-ph/0406039](#)]

$$|\lambda''_{11k}| \lesssim 10^{-7} \left( \frac{\tilde{m}}{100 \text{ GeV}} \right)^{5/2} \quad n - \bar{n}$$

$$|\lambda''_{312,313}| \lesssim O(10^{-2}), \quad (100 \text{ GeV} \lesssim \tilde{m} \lesssim 200 \text{ GeV}) \quad n - \bar{n}$$

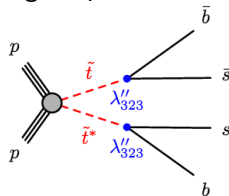
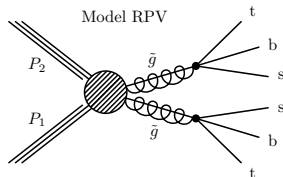
$$|\lambda''_{i23} \lambda''_{i13}|_{i=2,3} \lesssim 10^{-3} \left( \frac{m_{\tilde{u}_i}}{100 \text{ GeV}} \right) \quad K - \bar{K}$$

Matter bias, or suggesting a 3<sup>rd</sup> generation dominance?

Flavor symmetries automatically generate hierarchies. (see backup slides)

# RPV - LHC limits (so far)

No MET but large multiplicities, large  $H_T$ .



- gluino  $\rightarrow$  3j:

[CMS, 1208.2931], [ATLAS, 1502.05686]

$$m_{\tilde{g}} \gtrsim 1 \text{ TeV (model-independent)}$$

[Evans, Kats, Shih, Strassler, 1310.5758]  
[Graham, Rajendran, Saraswat, 1403.7197]

- squark  $\rightarrow$  2j: first limits on LSP squarks

$$m_{\tilde{t}} \gtrsim 350 - 400 \text{ GeV}$$

[CMS, 1412.7706], [ATLAS, ATLAS-CONF-2015-026]

GREAT! Finally limits on RPV squarks. Any blind spot?

# RPV - a step back to baryogenesis

$udd$  breaks baryon number:

$$\text{washout of } \Delta B \implies \lambda''_{\text{all}} < 10^{-7} \quad \text{for } T_{BG} \gtrsim \tilde{m}/10$$

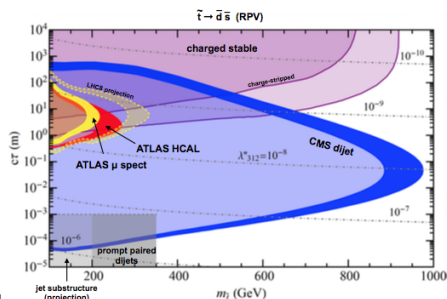
[Dreiner, Ross, Nuc1.Phys.B 1993]

Coincidentally,  $\lambda'' < 10^{-7} \implies$  displaced vertices at LHC!

[Barry, Graham, Rajendran, 1310.3853]

Limits are even stronger than prompt RPV! (see JP Chou's talk)

[Liu, Tweedie, 1503.05923]



$$m_{\tilde{\tau}} \gtrsim 800 - 1000 \text{ GeV}$$

[Csaki et al., 1505.00784]

[Zwane, 1505.03479]

## Large RPV

On the other hand, if  $\Delta B$  generated at  $T < \tilde{m}/10 \approx \mathcal{O}(10 - 100)$  GeV, no washout.

Many models use baryonic RPV for Baryogenesis:

Dimopoulos, Hall (1987), Cline, Raby (1991), Cui (2013), AM, Shin (2014), Arcadi, Covi, Nardecchia (2015)

For all models,  $X \rightarrow B, \bar{B}$

$$\varepsilon \equiv \frac{\Gamma_{X \rightarrow B} - \bar{\Gamma}_{X \rightarrow \bar{B}}}{2\Gamma} \propto \frac{|\lambda''|^2}{(16\pi^2)^\#} \Phi_{CP} f(\tilde{m}_i \dots)$$

$$10^{-10} \simeq \frac{n_B}{s} = \varepsilon \left( \frac{n_X}{s} \right)_{t=\frac{1}{\Gamma}} \implies \boxed{|\lambda''| \gtrsim \mathcal{O}(10^{-2} - 10^{-3})} \quad (*)$$

(\*loopholes: if X non-thermal. . . )

# A new channel at LHC

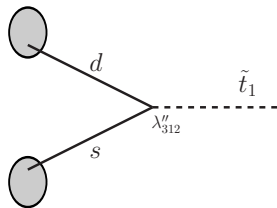
Experimental constraints and connection to baryogenesis *strongly suggests* that natural RPV SUSY involves large couplings (at least one):

resonant RPV production of squarks!

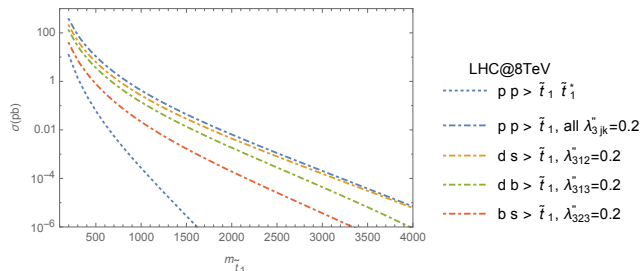
(will focus on 3rd gen. couplings for rest of talk)

$\lambda''_{3ij}$ ,  $3ij \equiv tds, tdb, tbs$

$$\hat{\sigma} = \frac{8\pi}{3} \frac{|\lambda''_{ijk}|^2}{m_{\tilde{t}}^2} \sin^2 \theta_{\tilde{t}} \delta(\hat{s} - m_{\tilde{t}}^2)$$



Large cross sections:  
(even *tbs*!)



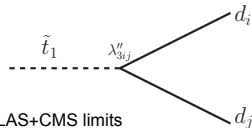
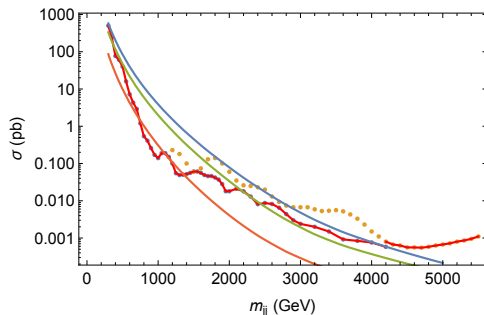
Disclaimer: old idea. [Dreiner,Ross, 1991]. There are (some) sensitivity studies for Tevatron [Berger,Harris,Sullivan, 1999] and LHC [Choudhury,Datta,Maity, 2011]. Also, CMS and ATLAS do a resonant  $\tilde{\nu}$  search (LRPV).

# Search strategies I: dijet resonance

$pp \rightarrow \tilde{t}_1 \rightarrow ds, db, bs$ : dijet resonance

[ATLAS-EXOT-2013-11,  $20.3\text{fb}^{-1}@8\text{TeV}$ ]

[CMS-EXO-12-059,  $19.7\text{fb}^{-1}@8\text{TeV}$ ]



— ATLAS+CMS limits

—  $d s \rightarrow \tilde{t}_1 \rightarrow d s, \lambda_{312}''=1$

—  $d b \rightarrow \tilde{t}_1 \rightarrow d b, \lambda_{313}''=1$

—  $b s \rightarrow \tilde{t}_1 \rightarrow b s, \lambda_{323}''=1$

No signal  $\Rightarrow$  Limits from 500 GeV to 5+ TeV.

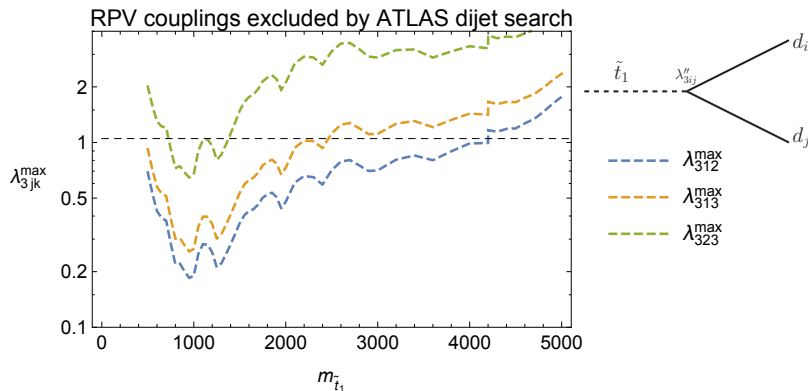


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Best limits around 1 TeV. Worst at low  $E$  (background), some searches did not even look. (CMS had specific trigger at end of Run I). (see backup slides)

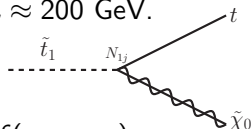
## Search strategies II: add a neutralino

A light higgsino is also favored by naturalness,  $\mu \approx 200$  GeV.

With  $\tilde{\chi}_0$  LSP, dijet signal suppressed,

$$Br(\tilde{t} \rightarrow jj) \approx \frac{\lambda''^2}{\lambda''^2 + (g^2 N_{11}^2 + \dots)} f(m_{\tilde{t}}, m_{\tilde{\chi}_0})$$

$$Br(\tilde{t} \rightarrow t\tilde{\chi}_0) \approx \frac{(g^2 N_{11}^2 + \dots)}{\lambda''^2 + (g^2 N_{11}^2 + \dots)} h(m_{\tilde{t}}, m_{\tilde{\chi}_0})$$



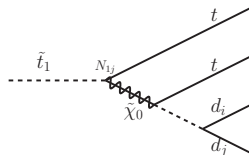
Dominant decay channel for  $\lambda'' \lesssim 0.1$ . Even for  $\lambda'' > 0.1$ ,  $\sigma \times Br_{t\tilde{\chi}_0} \approx \text{const.}$

Signatures? LSP is unstable:

- $c\tau > 10m \implies$  monotop? *unlikely*.
- $c\tau \ll 1m \implies$  prompt decays:

SS tops, constrained by CMS-SUS-13-013,

$$\sigma(pp \rightarrow tt, \bar{t}\bar{t}) < 720 \text{ fb}$$



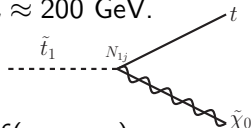
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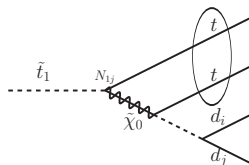


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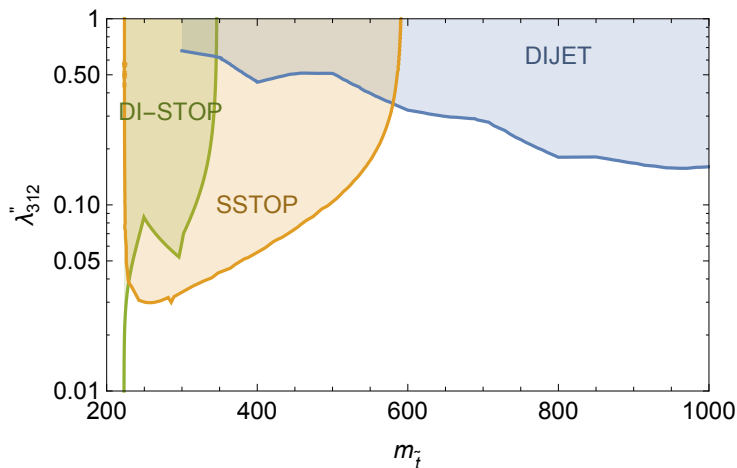
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# Combined New Limits below 1 TeV



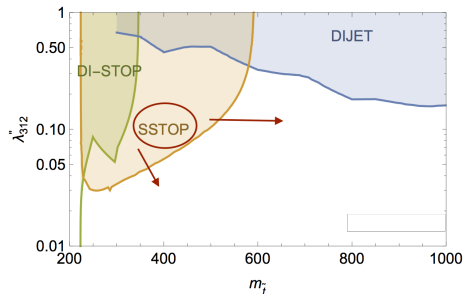
Here limits on  $\lambda''_{312}$  ( $tds$ ). (check out backup slides for  $\lambda''_{313}$ ,  $\lambda''_{323}$ .)

# Experimental outlook

How to close those gaps and improve limits/discover something:

(recommendations + new proposed searches in RED)

- SStop:  $\tilde{t} \rightarrow S S \ell + 2j$ : count extra jets, higher luminosity, more data.
- jj searches: more data, but do not forget low masses (trigger level?).
- Di-stop: more data, complementary  $4t + 4j$  search.



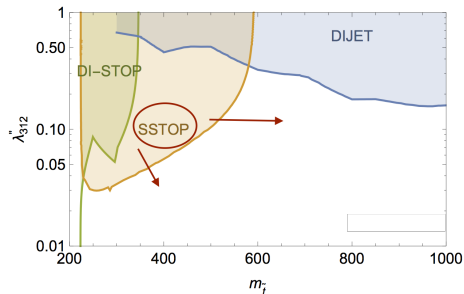
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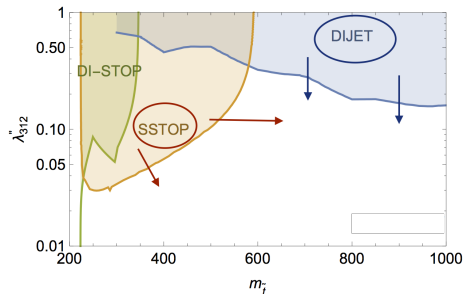
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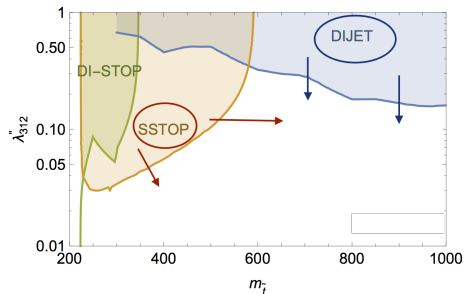
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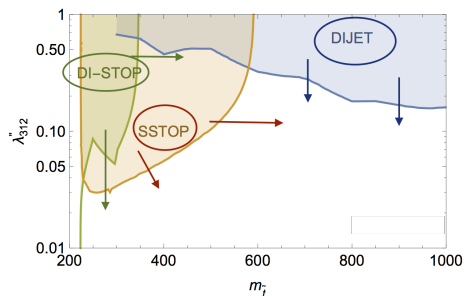


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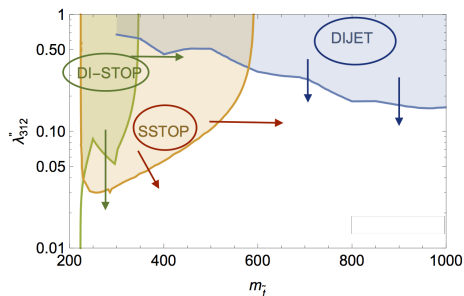
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# Conclusions

I have presented a series of **NEW** limits on RPV stops when those are produced through resonant scattering through the RPV interaction.

This was motivated by both null experimental searches for RPV below 1 TeV as well as the role of BNV interactions in baryogenesis.

Although mentioned before in the literature, it was overlooked so far in the analysis of LHC data. These are the **first direct limits** on individual RPV couplings, as opposed to flavor physics constraints.

At 13 TeV, cross section increases by a factor of  $\sim 10$ . We should keep looking at low mass range, interesting surprises can hide around the corner.

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## Extra slides

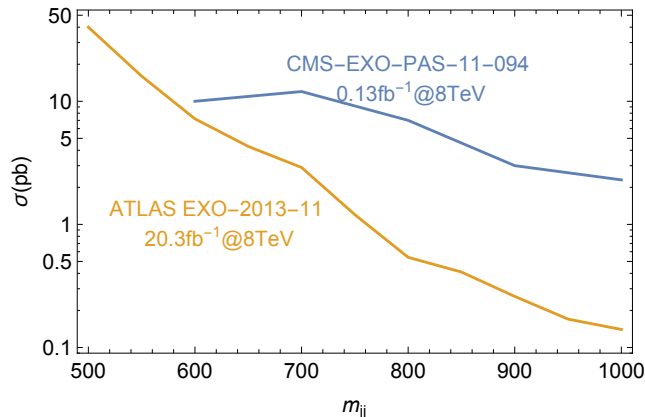
## Low-mass trigger: CMS-EXO-PAS-11-094

L1 trigger:  $H_T > 100$  GeV

HLT;  $H_T > 350$  GeV | ( $m_{jj} > 400$  GeV &  $|\Delta\eta_{jj}| < 2$ );

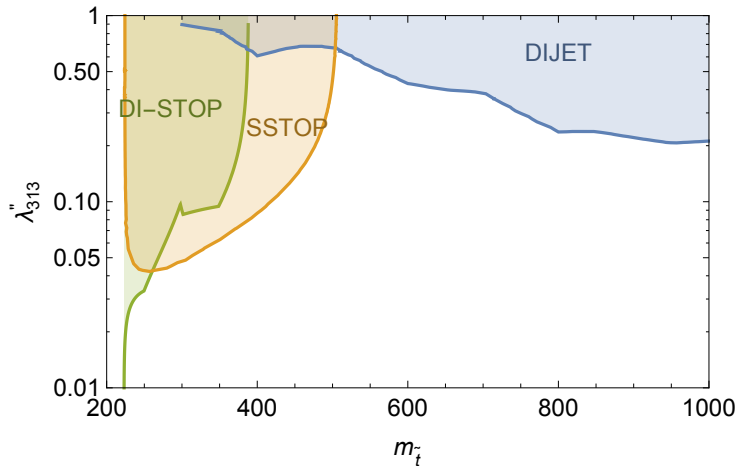
throw away most of event to reduce bandwidth.

Trigger active only at end of Run I: only  $0.13\text{fb}^{-1}$ ! Limits are quite good

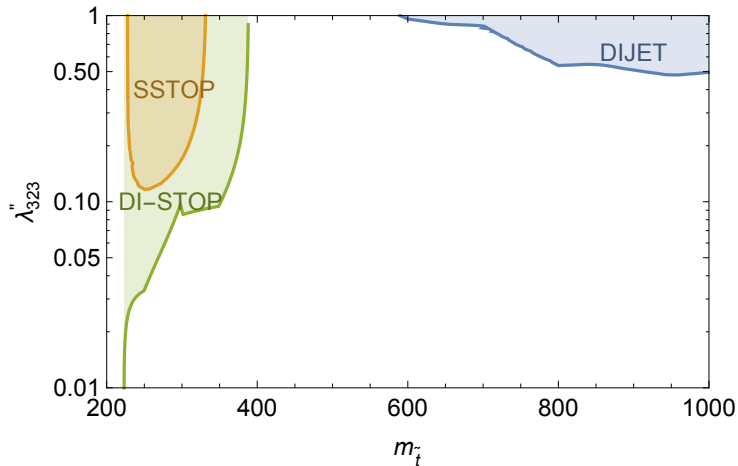


With  $20\text{fb}^{-1}$ , probe  
 $0.1\text{pb} \rightarrow \lambda'' \gtrsim 0.01!$

Limits on  $\lambda''_{313}$  (tdb)

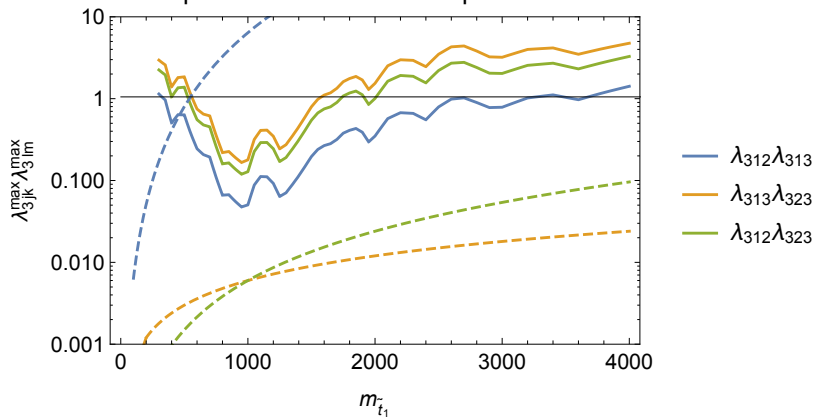


Limits on  $\lambda''_{323}$  (tbs)





# Comparison to $K-\bar{K}$ limits on products of $\lambda''$

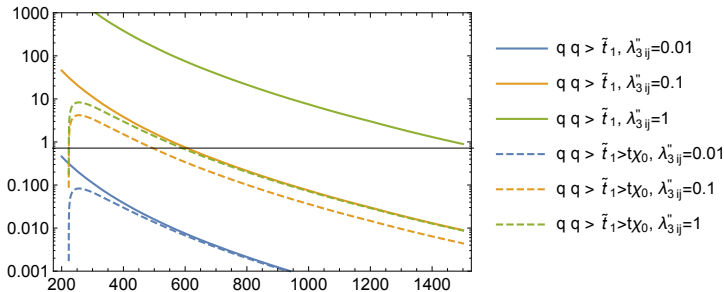


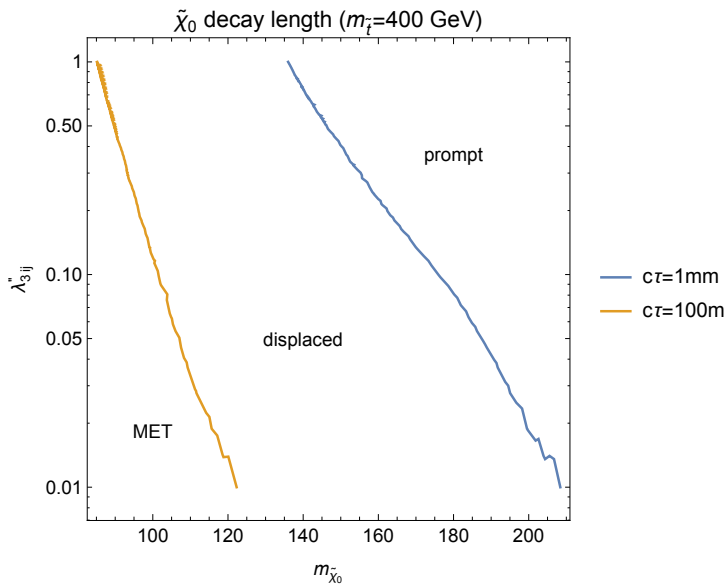
$pp \rightarrow \tilde{t} \rightarrow t\tilde{\chi}_0$  independent of  $\lambda''$  for large  $\lambda''$ :

$$\sigma \times Br(\tilde{t} \rightarrow t\tilde{\chi}_0) \sim \lambda''^2 \times \frac{(g^2 N_{11}^2 + \dots)}{\lambda''^2 + (g^2 N_{11}^2 + \dots)}$$

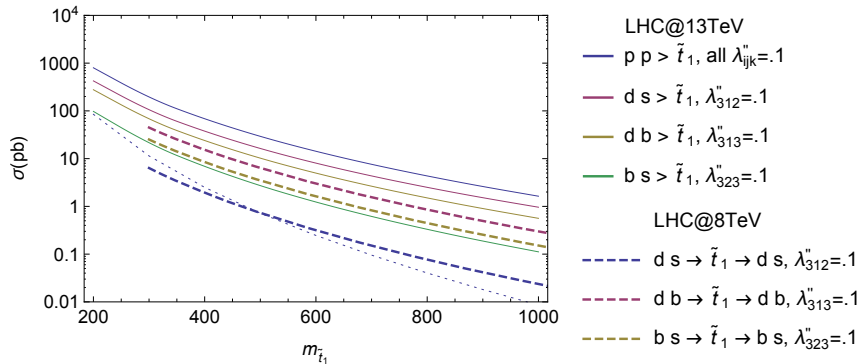
$$\simeq \begin{cases} (g^2 N_{11}^2 + \dots), & \lambda'' \gtrsim 0.1 \\ \lambda''^2 \times \mathcal{O}(1), & \lambda'' \lesssim 0.1 \end{cases}$$

312





# Cross sections at 13 TeV



# RPV - flavor symmetries

Flavor models also favor hierarchies with large 3<sup>rd</sup> generation couplings.

- horizontal flavor symmetries: Froggatt-Nielsen mechanism, known fermion masses and mixing *imply* hierarchies in RPV sector. [AM, 1305.2921]

$$\begin{pmatrix} \lambda''_{112} & \lambda''_{212} & \lambda''_{312} \\ \lambda''_{113} & \lambda''_{213} & \lambda''_{313} \\ \lambda''_{123} & \lambda''_{223} & \lambda''_{323} \end{pmatrix} = \lambda''_{323} \begin{pmatrix} 3 \times 10^{-5} & 3 \times 10^{-3} & 5 \times 10^{-2} \\ 10^{-4} & 10^{-2} & 2 \times 10^{-1} \\ 6 \times 10^{-4} & 5 \times 10^{-2} & 1 \end{pmatrix}$$

- MFV: assume  $Y_{ij}^f$  only flavor violation as spurion of SM flavor  $SU(3)^5$ .

$$\begin{pmatrix} \lambda''_{112} & \lambda''_{212} & \lambda''_{312} \\ \lambda''_{113} & \lambda''_{213} & \lambda''_{313} \\ \lambda''_{123} & \lambda''_{223} & \lambda''_{323} \end{pmatrix} = \frac{1}{2} \left( \frac{\tan \beta}{50} \right)^2 \begin{pmatrix} 10^{-8} & 10^{-4} & 2 \times 10^{-1} \\ 3 \times 10^{-5} & 5 \times 10^{-2} & 3 \times 10^{-1} \\ 2 \times 10^{-3} & 2 \times 10^{-1} & 1 \end{pmatrix}$$

[Nikolidakis, Smith, 0710.3129]  
[Csaki, Grossman, Heidenreich, 1111.1239]